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# THE STURM-LIOUVILLE PROBLEM WITH SINGULAR POTENTIAL AND THE EXTREMA OF THE FIRST EIGENVALUE

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ABSTRACT. We get the infima and suprema of the first eigenvalue of the problem

$$\begin{cases} -y'' + qy = \lambda y, \\ y'(0) - k_0^2 y(0) = 0, \\ y'(1) + k_1^2 y(1) = 0, \end{cases}$$

where q belongs to the set of constant-sign summable functions on [0, 1] such that

$$\int_{0}^{1} q \, dx = 1 \quad \text{or} \quad \int_{0}^{1} q \, dx = -1.$$

## 1. Introduction

**1.1.** Consider the Sturm–Liouville problem

$$-y'' + (q - \lambda)y = 0,$$
(1)
$$\int y'(0) - k_0^2 y(0) = 0.$$

$$\begin{cases} y'(0) - k_0 y(0) = 0, \\ y'(1) + k_1^2 y(1) = 0, \end{cases}$$
(2)

where the real coefficients  $k_0 \ge 0$  and  $k_1 \ge k_0$  are fixed, the solution y belongs to the space  $W_1^2[0, 1]$ , the equality (1) is considered as holding almost everywhere at [0, 1], and the potential  $q \in L_1[0, 1]$  is a constant-sign function such that one

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of the integral conditions holds:

$$\int_{0}^{1} q \, dx = 1 \quad \text{or} \quad \int_{0}^{1} q \, dx = -1. \tag{3}$$

The aim of this paper is to get the infima and suprema of the first eigenvalue of the problem (1)-(3).

**1.2.** The problem (1)–(3) is a partial case of the problem (1), (2) with  $q \in A_{\gamma}$  or  $-q \in A_{\gamma}$ , where  $\gamma \in \mathbb{R} \setminus \{0\}$  and

$$A_{\gamma} \rightleftharpoons \left\{ q \in L_1[0,1] : q(x) \ge 0 \text{ a.e. and } \int_0^1 q^{\gamma} dx = 1 \right\}.$$
(4)

Denote by  $\lambda_1(q)$  the minimal eigenvalue of the problem (1) or

$$-y'' - \lambda qy = 0 \tag{5}$$

with some self-adjoint boundary conditions. Consider for each  $\gamma \in \mathbb{R} \setminus \{0\}$  four values  $m_{\gamma}^{\pm} \rightleftharpoons \inf_{q \in A_{\gamma}} \lambda_1(\pm q)$  and  $M_{\gamma}^{\pm} \rightleftharpoons \sup_{q \in A_{\gamma}} \lambda_1(\pm q)$ . The estimates of  $m_{\gamma}^+$  and  $M_{\gamma}^+$  for the equation (5) with the Dirichlet boundary conditions were obtained in [1]. The analogous results about the Dirichlet problem for the equation (1) were obtained in [2], [3]. In [4] the problem (5), (2) was studied.

The values  $m_{\gamma}^+$  and  $M_{\gamma}^+$  for the problem (1), (2) with  $q \in A_{\gamma}$  were considered by one of the authors in [5] for all  $\gamma \neq 0$ . The most detailed and precise results were obtained for the case  $\gamma \neq 1$ .

The case  $\gamma = 1$  is in some kind special. In [3] and [5], for (1) with various boundary conditions, the precise results for  $M_1^+$  were obtained by the method quite different from used for  $\gamma \neq 1$ . In [5] for  $m_1^+$  only inequality  $m_1^+ \geq 1/4$ was obtained. In [3] for  $m_1^-$  it was proved that this infimum is attained at the non-summable potential  $q^* = -\delta_{1/2}$ .

In this paper we extend the class of considered potentials from  $L_1[0, 1]$  to the space  $W_2^{-1}[0, 1]$  (see [6] and 2.1 later). The space  $W_2^{-1}[0, 1]$ , in particular, contains a Dirac delta function  $\delta_{\zeta}$  with support located at an arbitrary point  $\zeta \in [0, 1]$ . This generalization of the problem lets us to get the precise description of  $M_1^-$  and  $m_1^{\pm}$  and to prove that they are attained at the potentials from the extended class.

**1.3.** The main results of the paper are the following four theorems:

**1.3.1. THEOREM.** By definition, put

$$\alpha_{\mu} \rightleftharpoons \frac{1}{\sqrt{\mu}} \arctan \frac{k_0^2}{\sqrt{\mu}}, \quad \beta_{\mu} \rightleftharpoons \frac{1}{\sqrt{\mu}} \arctan \frac{k_1^2}{\sqrt{\mu}}.$$
(6)

Then  $M_1^+$  is a unique solution to the equation

$$1 - \alpha_{\mu} - \beta_{\mu} = \mu^{-1} \tag{7}$$

and is attained at the potential  $q^* \in L_1[0,1]$  such that

$$q^*(x) = \begin{cases} M_1^+ & \text{for } x \in \left[\alpha_{M_1^+}, 1 - \beta_{M_1^+}\right], \\ 0 & \text{otherwise.} \end{cases}$$

**1.3.2. THEOREM.** If  $k_0^2 + k_1^2 \leq 1$ , then  $M_1^- = k_0^2 + k_1^2 - 1$  and is attained at the potential  $a^* \rightarrow b^2 \delta = b^2 \delta = (1 - b^2 - b^2)$ 

$$q^* \rightleftharpoons -k_0^2 \boldsymbol{\delta}_0 - k_1^2 \boldsymbol{\delta}_1 - \left(1 - k_0^2 - k_1^2\right).$$

If  $k_0^2 + k_1^2 \ge 1$  and  $k_1^2 - k_0^2 \le 1$ , then  $M_1^-$  is the minimal eigenvalue of the problem

$$-y'' = \lambda y, \tag{8}$$

$$2y'(0) - \left(k_0^2 + k_1^2 - 1\right)y(0) = 2y'(1) + \left(k_0^2 + k_1^2 - 1\right)y(1) = 0$$
(9)

and is attained at the potential

$$q^* \cong -(1+k_0^2-k_1^2)\boldsymbol{\delta}_0/2 - (1-k_0^2+k_1^2)\boldsymbol{\delta}_1/2.$$

If  $k_1^2 - k_0^2 \ge 1$ , then  $M_1^-$  is the minimal eigenvalue of the problem (8) with

$$y'(0) - k_0^2 y(0) = y'(1) + (k_1^2 - 1)y(1) = 0$$
(10)

and is attained at the potential  $q^* \rightleftharpoons - \boldsymbol{\delta}_1$ .

**1.3.3. THEOREM.**  $m_1^+$  is the minimal eigenvalue of the problem (8) with

$${}'(0) - k_0^2 y(0) = y'(1) + (k_1^2 + 1)y(1) = 0$$
(11)

 $y'(0) - k_0^2 y(0) = y'(1)$ and is attained at the potential  $q^* \rightleftharpoons \boldsymbol{\delta}_1$ .

**1.3.4. Theorem.** If for some  $\mu \geq -k_0^4$  and some  $\zeta \in (0,1)$  the problem

$$-y'' = \mu y$$
 at  $(0, \zeta) \cup (\zeta, 1),$  (12)

$$y'(0) - k_0^2 y(0) = 2y'(\zeta - 0) - y(\zeta)$$
  
= 2y'(\zeta + 0) + y(\zeta) = y'(1) + k\_1^2 y(1) = 0 (13)

has a continuous positive solution, then  $m_1^- = \mu$  and  $m_1^-$  is attained at the potential  $q^* \rightleftharpoons -\delta_{\zeta}$ . Otherwise  $m_1^-$  is the minimal eigenvalue of the problem (8) with

$$y'(0) - (k_0^2 - 1)y(0) = y'(1) + k_1^2 y(1) = 0$$

and is attained at the potential  $q^* \rightleftharpoons -\boldsymbol{\delta}_0$ .

Some additional remarks on solvability of the boundary problem (12), (13) will be given in the subsection 3.6.

**1.4.** Let us give some examples that illustrate the theorems from the previous subsection. In the case  $k_0 = k_1 = 0$  we get  $m_1^+ = \lambda_1(\boldsymbol{\delta}_1) = 0.740174(\pm 10^{-6})$ . In the case  $k_0^2 = k_1^2 > 1/2$  we get  $m_1^- = \lambda_1(-\boldsymbol{\delta}_{1/2})$ . In the case  $k_0^2 = k_1^2 = 1/2$  we have  $m_1^- = \lambda_1(-\boldsymbol{\delta}_{\zeta}) = -1/4$  for any  $\zeta \in [0, 1]$ . In the case  $k_0^2 = k_1^2 < 1/2$  we have  $m_1^- = \lambda_1(-\boldsymbol{\delta}_0)$ .

## **2.** The set $\Gamma_1$ and related topics

2.1. We suppose that all considered functional spaces are real.

By  $W_2^{-1}[0,1]$  denote the Hilbert space that is a completion of  $L_2[0,1]$  in the norm

$$\|y\|_{W_2^{-1}[0,1]} \rightleftharpoons \sup_{\|z\|_{W_2^{1}[0,1]}=1} \int_0^1 yz \, dx.$$

When  $y \in W_2^{-1}[0,1]$ , by  $\int_0^1 yz \, dx$  we sometimes denote the result

$$\langle y, z \rangle \rightleftharpoons \lim_{n \to \infty} \int_{0}^{1} y_n z \, dx, \quad \text{where} \quad y = \lim_{n \to \infty} y_n, \ y_n \in L_2[0, 1],$$

of applying the linear functional y to the function  $z \in W_2^1[0,1]$ .

For any fixed  $q \in L_1[0,1]$  and  $\lambda \in \mathbb{R}$  the map taking each  $y \in W_1^2[0,1]$  satisfying (2) to

$$-y'' + (q - \lambda)y \in L_1[0, 1]$$

can be extended by continuity to the bounded operator  $T_q(\lambda) : W_2^1[0,1] \to W_2^{-1}[0,1]$ . Using integration by part, we get

$$(\forall y, z \in W_2^1[0, 1]) \quad \left\langle T_q(\lambda)y, z \right\rangle$$
  
=  $\int_0^1 \left[ y'z' + (q - \lambda)yz \right] dx + k_0^2 y(0)z(0) + k_1^2 y(1)z(1).$ (14)

Consider the linear operator pencil<sup>1</sup>  $T_q : \mathbb{R} \to \mathcal{B}(W_2^1[0,1], W_2^{-1}[0,1])$  that takes any  $\lambda \in \mathbb{R}$  to the operator  $T_q(\lambda)$  described by (14). The spectral problem for  $T_q$  may be considered as a reformulation (or as a generalization in case when  $q \in W_2^{-1}[0,1]$  is not summable) of the boundary value problem (1), (2). We can do this due to the following two facts.

<sup>&</sup>lt;sup>1</sup>A linear operator pencil L is an operator-valued function such that  $L(\lambda) = A + \lambda B$ , where  $\lambda \in \mathbb{R}$ , A and B are some operators not depending on  $\lambda$ .

**2.1.1.** For all  $q \in L_1[0,1]$  and  $\lambda \in \mathbb{R}$  the function  $y \in W_2^1[0,1]$  belongs to the kernel of the operator  $T_q(\lambda)$ , if and only if  $y \in W_1^2[0,1]$  and y is a solution of the problem (1), (2).

Proof. It directly follows from the definition of the operator  $T_q(\lambda)$  that for any solution  $y \in W_1^2[0,1]$  of the problem (1), (2) the equality  $T_q(\lambda)y = 0$  holds.

Let us prove the converse. Consider some  $y \in \ker T_q(\lambda)$ , and put

$$w(x) \rightleftharpoons y'(x) - \int_{0}^{x} (q - \lambda)y \, dt.$$
(15)

For any  $z \in \overset{\circ}{W}{}_{2}^{1}[0,1]$ , using (14), we have

$$0 = \left\langle T_q(\lambda)y, z \right\rangle = \int_0^1 w z' \, dx. \tag{16}$$

Since the set of the derivatives of all functions  $z \in \mathring{W}_2^1[0,1]$  is an orthogonal complement in  $L_2[0,1]$  of the set of all constants, from (16) it follows that the function  $w \in L_2[0,1]$  is constant. Combining this with (15), we get that the function y' is absolutely continuous and its generalized derivative equals  $(q - \lambda)y$ . Now, using (14), we see that for any  $z \in W_2^1[0,1]$  we get

$$0 = \langle T_q(\lambda)y, z \rangle = \left[ -y'(0) + k_0^2 y(0) \right] z(0) + \left[ y'(1) + k_1^2 y(1) \right] z(1),$$
  
tisfies the conditions (2).

so y satisfies the conditions (2).

**2.1.2.** For any  $q \in W_2^{-1}[0,1]$  the spectrum of the linear operator pencil  $T_q$  is purely discrete, simple and bounded from below.

Proof. Note that for any  $y \in W_2^1[0,1]$  we have

$$\|y^2\|_{W_2^1[0,1]} \le \sup_{x \in [0,1]} |y(x)| \cdot \sqrt{\int_0^1 \left[y^2 + 4(y')^2\right] dx} \le 2\|y\|_{C[0,1]} \cdot \|y\|_{W_2^1[0,1]},$$

then, by the embedding theorem, we get

$$\|y^2\|_{W_2^1[0,1]} \le C \, \|y\|_{W_2^1[0,1]}^2,\tag{17}$$

where C is some constant.

Since C[0,1] is densely embedded in  $W_2^{-1}[0,1]$ , for any  $\varepsilon \in (0,1)$  there exists a function  $\tilde{q} \in C[0,1]$  such that

$$\|\tilde{q} - q\|_{W_2^{-1}[0,1]} \le \varepsilon/C.$$

Using this and the inequality (17), for any  $y \in W_2^1[0,1]$  we get

$$\left| \int_{0}^{1} (\tilde{q} - q) y^{2} dx \right| \leq \|\tilde{q} - q\|_{W_{2}^{-1}[0,1]} \cdot \|y^{2}\|_{W_{2}^{1}[0,1]} \leq \varepsilon \|y\|_{W_{2}^{1}[0,1]}^{2}.$$
(18)

Further, for any  $\kappa > \|\tilde{q}\|_{C[0,1]} + 1$  we have  $\int_0^1 \tilde{q}y^2 dx \ge (1-\kappa) \int_0^1 y^2 dx$ . Combining this with (14) and (18), we obtain

$$IT_q(-\kappa) \ge 1 - \varepsilon, \tag{19}$$

where by  $I: W_2^{-1}[0,1] \to W_2^{1}[0,1]$  we denote an isometry that satisfies

$$(\forall y \in W_2^{-1}[0,1]) (\forall z \in W_2^1[0,1]) \qquad \langle Iy, z \rangle_{W_2^1[0,1]} = \langle y, z \rangle.$$

The existence and uniqueness of this isometry follows from the Riesz theorem about the representation of a functional in a Hilbert space  $[7, \S 30, \S 99]$ .

From the estimate (19) it follows [7, § 104] that the operator  $S \rightleftharpoons IT_q(-\kappa)$  is boundedly invertible. Taking into account (14), we have  $IT_q(\lambda) \equiv S - (\lambda + \kappa)J^*J$ , where  $J: W_2^1[0, 1] \to L_2[0, 1]$  is the embedding operator. So for any  $\lambda \in \mathbb{R}$  the existence of a bounded inverse of the operator  $T_q(\lambda)$  is equivalent to the existence of a bounded inverse of the operator  $1 - (\lambda + \kappa)S^{-1/2}J^*JS^{-1/2}$ . Since J is compact, it follows that the spectrum of  $T_q$  is purely discrete, semi-simple and bounded from below.

The spectrum of the pencil  $T_q$  is simple since (see [6], [8, Propositions 2, 10]) for any  $\lambda \in \mathbb{R}$  the kernel of the operator  $T_q(\lambda)$  is formed by the first components  $Y_1$  of the solutions to the boundary value problem

$$\begin{pmatrix} Y_1 \\ Y_2 \end{pmatrix}' = \begin{pmatrix} u & 1 \\ -u^2 & -u \end{pmatrix} \begin{pmatrix} Y_1 \\ Y_2 \end{pmatrix},$$
(20)

$$Y_2(0) - k_0^2 Y_1(0) = Y_2(1) + \left[k_1^2 + \omega\right] Y_1(1) = 0.$$
(21)

Here  $u \in L_2[0,1]$  and  $\omega \in \mathbb{R}$  are taken from the representation

$$(\forall y \in W_2^1[0,1]) \qquad \int_0^1 (q-\lambda)y \, dx = -\int_0^1 uy' \, dx + \omega \, y(1) \tag{22}$$

of the potential  $q \in W_2^{-1}[0, 1]$ .

**2.2.** For the eigenvalues

$$\lambda_1(q) < \lambda_2(q) < \cdots < \lambda_n(q) < \cdots$$

of the pencil  $T_q$  we have the following propositions.

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**2.2.1.** (See [8, Proposition 10].) For any  $n \ge 1$ ,  $q \in W_2^{-1}[0,1]$  and  $\lambda \in \mathbb{R}$  the inequality  $\lambda > \lambda_n(q)$  is equivalent to the existence of n-dimensional subspace  $\mathfrak{N} \subset W_2^1[0,1]$  that satisfies

$$(\forall y \in \mathfrak{N} \setminus \{0\}) \qquad \langle T_q(\lambda)y, y \rangle < 0.$$

**2.2.2.** For any  $n \ge 1$  the function  $\lambda_n \colon W_2^{-1}[0,1] \to \mathbb{R}$  is continuous.

Proof. Consider some  $q \in W_2^{-1}[0,1]$  and  $\varepsilon \in (0,1/2)$ . For any  $y \in W_2^1[0,1]$ ,  $\lambda \in \mathbb{R}$  and  $\tilde{q} \in W_2^{-1}[0,1]$  such that  $\|\tilde{q} - q\|_{W_2^{-1}[0,1]} < \varepsilon/C$ , where C is the same as in (17), we get

$$\begin{split} \left\langle T_{\bar{q}}(\lambda)y,y\right\rangle &\geq \left\langle T_{q}(\lambda)y,y\right\rangle - \varepsilon \,\|y\|_{W_{2}^{1}[0,1]}^{2}\\ &\geq \left\langle T_{q}(\lambda)y,y\right\rangle - \varepsilon \,\|y\|_{W_{2}^{1}[0,1]}^{2}\\ &\quad -\varepsilon \cdot \left\langle T_{2q}(\lambda_{1}(2q))y,y\right\rangle - \varepsilon k_{0}^{2}y^{2}(0) - \varepsilon k_{1}^{2}y^{2}(1)\\ &= (1-2\varepsilon) \cdot \left\langle T_{q}\left(\frac{\lambda + \varepsilon \cdot [1-\lambda_{1}(2q)]}{1-2\varepsilon}\right)y,y\right\rangle. \end{split}$$

Consequently, from the variational principle 2.2.1 it follows that any  $\lambda > \lambda_n(\tilde{q})$  satisfies

$$\frac{\lambda + \varepsilon \cdot [1 - \lambda_1(2q)]}{1 - 2\varepsilon} > \lambda_n(q).$$

Since we can choose  $\lambda$  arbitrarily close to  $\lambda_n(\tilde{q})$ , we have

$$\lambda_n(\tilde{q}) \ge (1-2\varepsilon)\,\lambda_n(q) - \varepsilon \cdot [1-\lambda_1(2q)].$$

By the same method we get

$$\lambda_n(\tilde{q}) \leq (1+2\varepsilon)\,\lambda_n(q) + \varepsilon \cdot [1-\lambda_1(2q)].$$

**2.3.** Let  $\Gamma_1$  be the closure in  $W_2^{-1}[0,1]$  of the set  $A_1$  defined by (4). Put by definition

$$\Lambda(X) \rightleftharpoons \big\{ \lambda \in \mathbb{R} : (\exists q \in X) \quad \lambda = \lambda_1(q) \big\},\$$

where  $X \subseteq W_2^{-1}[0,1]$  is some set of generalized functions. The set  $\Lambda(X)$  is formed by all the possible values of  $\lambda_1(q)$  for all  $q \in X$ . By -X we, as usually, denote the set

$$\{q \in W_2^{-1}[0,1] : (\exists r \in X) \quad q = -r\}.$$

**2.3.1.** Suppose X is a dense subset of  $\Gamma_1$ , then the closures of  $\Lambda(\pm X)$  and  $\Lambda(\pm\Gamma_1)$  coincide.

**2.3.2.** The extrema  $m_1^{\pm} \rightleftharpoons \inf \Lambda(\pm A_1)$  and  $M_1^{\pm} \rightleftharpoons \sup \Lambda(\pm A_1)$ , defined in 1.2, satisfy the equalities  $m_1^{\pm} \doteq \inf \Lambda(\pm \Gamma_1)$  and  $M_1^{\pm} = \sup \Lambda(\pm \Gamma_1)$ .

The proposition 2.3.1 immediately follows from 2.2.2. The proposition 2.3.2 immediately follows from 2.3.1.

**2.3.3.** The set  $\Gamma_1$  consists of all non-negative<sup>2</sup> distributions  $q \in W_2^{-1}[0,1]$  such that  $\int_0^1 q \, dx = 1$ .

Proof. Since for any  $q \in \Gamma_1$  there exists a sequence of functions from  $A_1$  such that its limit equals q, it follows that the generalized function q is non-negative and satisfies  $\int_0^1 q \, dx = 1$ .

Let us prove the converse. Suppose  $q \in W_2^{-1}[0, 1]$  is a non-negative generalized function and satisfies  $\int_0^1 q \, dx = 1$ . Then (see [6], [8, § 2.3]) there exists a function  $u \in L_2[0, 1]$  such that

$$(\forall y \in W_2^1[0,1])$$
  $\int_0^1 qy \, dx = -\int_0^1 uy' \, dx + y(1).$  (23)

Put by definition

$$\Pi_{\gamma,\eta,\theta}(x) \rightleftharpoons \begin{cases} \frac{x-\gamma}{\eta-\gamma} & \text{for } x \in [\gamma,\eta];\\ \frac{\theta-x}{\theta-\eta} & \text{for } x \in [\eta,\theta],\\ 0 & \text{otherwise} \end{cases}$$

for any reals  $\gamma < \eta < \theta$ . Suppose 0 < a < b < c < d < 1. Substituting the functions  $\Pi_{-1,0,a} + \Pi_{0,a,b}$ ,  $\Pi_{a,b,c} + \Pi_{b,c,d}$  and  $\Pi_{c,d,1} + \Pi_{d,1,2}$  for y in (23), we get

$$0 \le \frac{1}{b-a} \int_{a}^{b} u \, dx \le \frac{1}{d-c} \int_{c}^{d} u \, dx \le 1.$$

From these inequalities it follows that the function  $u \in L_2[0,1]$  is non-decreasing and satisfies vrai  $\inf_{x \in [0,1]} u(x) \ge 0$  and vrai  $\sup_{x \in [0,1]} u(x) \le 1$ .

Since there exists a sequence  $\{u_n\}_{n=0}^{\infty}$  of non-decreasing piecewise linear functions such that  $u_n(0) = 0$ ,  $u_n(1) = 1$  and  $u = \lim_{n \to \infty} u_n$ , it follows that  $q = \lim_{n \to \infty} u'_n$ , where  $u'_n \in A_1$ .

**2.4.** Consider the function F implicitly defined by the equation

$$\lambda_1 \big( F(\mu, \zeta) \boldsymbol{\delta}_{\zeta} \big) = \mu, \tag{24}$$

where  $\mu \in \mathbb{R}$  and  $\zeta \in [0, 1]$ . The following three propositions give us some information about this function.

<sup>&</sup>lt;sup>2</sup>The generalized function  $q \in W_2^{-1}[0, 1]$  is called non-negative if for any non-negative function  $y \in W_2^1[0, 1]$  the inequality  $\langle q, y \rangle \ge 0$  holds.

**2.4.1.** For any  $\zeta \in [0,1]$  the function  $F(\cdot,\zeta)$  is single-valued, strictly increasing, and its domain is the interval  $(-\infty, f^+)$  with some  $f^+ > 0$ .

Proof. For any  $a \in \mathbb{R}$  there exists [8, Proposition 11] a positive eigenfunction  $y \in \ker T_{a\delta_{\zeta}}(\mu)$  corresponding to the eigenvalue  $\mu \rightleftharpoons \lambda_1(a\delta_{\zeta})$ , so for any b < a we have

$$\langle T_{b\boldsymbol{\delta}_{\zeta}}(\mu)y, y \rangle = \langle T_{a\boldsymbol{\delta}_{\zeta}}(\mu)y, y \rangle + (b-a) \cdot y^{2}(\zeta) < 0.$$

Using 2.2.1, we now get  $\lambda_1(b\delta_{\zeta}) < \mu$ . So the function  $F(\cdot, \zeta)$  is the inverse of the strictly increasing and, according to 2.2.2, continuous map  $a \mapsto \lambda_1(a\delta_{\zeta})$ . Therefore, the function  $F(\cdot, \zeta)$  is single-valued and strictly increasing.

Further, for any  $a \in \mathbb{R}$  from the equality

$$\left\langle T_{a\boldsymbol{\delta}_{\zeta}}\left(a+k_{0}^{2}+k_{1}^{2}\right)1,1\right\rangle = a-\left(a+k_{0}^{2}+k_{1}^{2}\right)+k_{0}^{2}+k_{1}^{2}=0$$

and the proposition 2.2.1 it follows that  $\lambda_1(a\delta_{\zeta}) \leq a + k_0^2 + k_1^2$ . Therefore, the domain of  $F(\cdot, \zeta)$  is unbounded from below. Also for any a > 0 we have  $\lambda_1(a\delta_{\zeta}) > 0$ , so the right bound of dom  $F(\cdot, \zeta)$  is positive.

## **2.4.2.** The function F is continuous.

Proof. Consider an arbitrary point  $(\mu_0, \zeta_0) \in \text{dom } F$  and suppose  $a^{\pm}$  satisfy  $a^- < F(\mu_0, \zeta_0) < a^+$ . For any point  $(\mu, \zeta) \in \mathbb{R} \times [0, 1]$  sufficiently close to  $(\mu_0, \zeta_0)$  from 2.4.1 and 2.2.2 we obtain the inequalities  $\lambda_1(a^-\delta_{\zeta}) < \mu < \lambda_1(a^+\delta_{\zeta})$ . Hence there exists  $a \in (a^-, a^+)$  such that  $\mu = \lambda_1(a\delta_{\zeta})$ , so for the point  $(\mu, \zeta)$  the equation (24) has a solution  $F(\mu, \zeta) = a$ .

**2.4.3.** A point  $(\mu, \zeta) \in (0, +\infty) \times [0, 1]$  belongs to domain of the function F if and only if the following conditions hold:

$$\sqrt{\mu} \cdot (\zeta - \alpha_{\mu}) \in (-\pi/2, \pi/2), \quad \sqrt{\mu} \cdot (1 - \beta_{\mu} - \zeta) \in (-\pi/2, \pi/2),$$
 (25)

where  $\alpha_{\mu}$  and  $\beta_{\mu}$  are defined by (6). In this case the equality

$$F(\mu,\zeta) = \sqrt{\mu} \cdot \left\{ \tan\left[\sqrt{\mu} \cdot (\zeta - \alpha_{\mu})\right] + \tan\left[\sqrt{\mu} \cdot (1 - \beta_{\mu} - \zeta)\right] \right\}$$
(26)

holds.

For any  $\zeta \in [0,1]$  the equality

$$F(0,\zeta) = -\frac{k_0^2}{1+k_0^2\zeta} - \frac{k_1^2}{1+k_1^2(1-\zeta)}$$
(27)

holds.

For any  $\mu < 0$  and  $\zeta \in [0, 1]$  the equality

$$F(\mu,\zeta) = -\sqrt{|\mu|} \cdot \left\{ G\left(\sqrt{|\mu|}, k_0^2, \zeta\right) + G\left(\sqrt{|\mu|}, k_1^2, 1-\zeta\right) \right\},$$
(28)

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where

$$G(\nu, \kappa, x) \rightleftharpoons \begin{cases} \tanh\left(\nu x + \ln\sqrt{\frac{\nu + \kappa}{\nu - \kappa}}\right) & \text{for } \nu > \kappa, \\ 1 & \text{for } \nu = \kappa, \\ \coth\left(\nu x + \ln\sqrt{\frac{\kappa + \nu}{\kappa - \nu}}\right) & \text{for } \nu < \kappa, \end{cases}$$

holds.

Proof. Consider  $\mu \in \mathbb{R}$  and  $\zeta \in (0, 1)$  such that  $(\mu, \zeta) \in \text{dom } F$ . According to (20)–(22), the equality  $T_q(\mu)y = 0$ , where  $q \rightleftharpoons F(\mu, \zeta)\delta_{\zeta}$ , is equivalent to the boundary problem

$$-y'' = \mu y$$
 at  $(0, \zeta) \cup (\zeta, 1)$ , (29)

$$y'(\zeta + 0) - y'(\zeta - 0) = F(\mu, \zeta)y(\zeta),$$
(30)

$$y'(0) - k_0^2 y(0) = y'(1) + k_1^2 y(1) = 0.$$
(31)

From [8, Proposition 11] and (24) it follows that any non-trivial solution to the problem (29)-(31) is constant-sign.

In the case  $\mu > 0$  any solution to the problem (29), (31) has the form

$$y(x) = \begin{cases} A \cdot \cos\left[\sqrt{\mu} \cdot (1 - \beta_{\mu} - \zeta)\right] \cdot \cos\left[\sqrt{\mu} \cdot (x - \alpha_{\mu})\right] & \text{for } x < \zeta, \\ A \cdot \cos\left[\sqrt{\mu} \cdot (1 - \beta_{\mu} - x)\right] \cdot \cos\left[\sqrt{\mu} \cdot (\zeta - \alpha_{\mu})\right] & \text{for } x > \zeta, \end{cases}$$
(32)

where A is some constant. This function is constant-sign if and only if the conditions (25) hold. Using (30), we now get (26). The values  $\zeta \in \{0, 1\}$  are finally included in the consideration using the propositions 2.4.2 and 2.2.2.

The cases  $\mu = 0$  and  $\mu < 0$  are considered on the base of (29)–(31) by analogous way using the solution

$$y(x) = \begin{cases} A \cdot [1 + k_1^2 (1 - \zeta)] \cdot [1 + k_0^2 x] & \text{for } x < \zeta, \\ A \cdot [1 + k_1^2 (1 - x)] \cdot [1 + k_0^2 \zeta] & \text{for } x > \zeta \end{cases}$$
(33)

in the case  $\mu = 0$ , and the solution

$$y(x) = \begin{cases} A \cdot g\left(\sqrt{|\mu|}, k_1^2, 1-\zeta\right) \cdot g\left(\sqrt{|\mu|}, k_0^2, x\right) & \text{for } x < \zeta, \\ A \cdot g\left(\sqrt{|\mu|}, k_1^2, 1-x\right) \cdot g\left(\sqrt{|\mu|}, k_0^2, \zeta\right) & \text{for } x > \zeta, \end{cases}$$
(34)

where

$$g(\nu, \kappa, x) \rightleftharpoons \begin{cases} \cosh\left(\nu x + \ln\sqrt{\frac{\nu + \kappa}{\nu - \kappa}}\right) & \text{for } \nu > \kappa, \\ e^{\nu x} & \text{for } \nu = \kappa, \\ \sinh\left(\nu x + \ln\sqrt{\frac{\kappa + \nu}{\kappa - \nu}}\right) & \text{for } \nu < \kappa, \end{cases}$$

in the case  $\mu < 0$ .

## 3. Proofs of the main results

**3.1.** In this section we prove Theorems 1.3.1–1.3.4. We use the notation

$$\Omega^+(y) \rightleftharpoons \left\{ x \in [0,1] : y(x) = \sup_{t \in [0,1]} y(t) \right\},$$
$$\Omega^-(y) \rightleftharpoons \left\{ x \in [0,1] : y(x) = \inf_{t \in [0,1]} y(t) \right\},$$

where  $y \in W_2^1[0, 1]$  is an arbitrary positive function. Also we take into account proposition 2.3.2.

**3.2. Proof of Theorem 1.3.1.** Consider some potential  $q^* \in \Gamma_1$ , and some positive eigenfunction  $y \in \ker T_{q^*}(\lambda_1(q^*))$ . Suppose that the support of the generalized function  $q^*$  is a subset of  $\Omega^+(y)$ . Then for any  $q \in \Gamma_1$  we, using 2.3.3, have

$$0 = \left\langle T_{q^*}(\lambda_1(q^*))y, y \right\rangle$$
  
=  $\int_0^1 [(y')^2 - \lambda_1(q^*)y^2] dx + \sup_{x \in [0,1]} y^2(x) + k_0^2 y^2(0) + k_1^2 y^2(1)$   
 $\ge \int_0^1 [(y')^2 + (q - \lambda_1(q^*))y^2] dx + k_0^2 y^2(0) + k_1^2 y^2(1),$ 

hence

$$\left\langle T_q(\lambda_1(q^*))y, y \right\rangle \le 0.$$

It follows that  $\lambda_1(q) \leq \lambda_1(q^*)$ , therefore  $\lambda_1(q^*) = M_1^+$ . Thus we have proved that  $M_1^+$  is attained at any potential  $q^*$  such that  $\sup q^* \subseteq \Omega^+(y)$ .

Suppose that  $\Omega^+(y) = [\tau_0, \tau_1]$ , where  $\tau_0 \neq \tau_1$ . Also suppose that the potential  $q^*$  is summable and has the form

$$q^*(x) = \begin{cases} \mu & \text{for } x \in [\tau_0, \tau_1], \\ 0 & \text{otherwise,} \end{cases}$$

where  $\mu$  is some positive constant. Since y''(x) = 0 for all  $x \in (\tau_0, \tau_1)$ , it follows that  $\mu = \lambda_1(q^*)$ . Therefore, the eigenfunction y has the form

$$y(x) = \begin{cases} A \cdot \cos\left[\sqrt{\mu} \cdot (x - \alpha_{\mu})\right] & \text{for } x < \tau_0, \\ B & \text{for } x \in [\tau_0, \tau_1], \\ C \cdot \cos\left[\sqrt{\mu} \cdot (1 - \beta_{\mu} - x)\right] & \text{for } x > \tau_1, \end{cases}$$

where A, B and C are some positive constants, and  $\alpha_{\mu}$ ,  $\beta_{\mu}$  are defined by (6). From the continuity of y' it follows that  $\tau_0 = \alpha_{\mu}$  and  $\tau_1 = 1 - \beta_{\mu}$ , hence A = B = C. Finally, from the condition  $\int_0^1 q^* dx = 1$  we have the equation (7).

To conclude the proof, it remains to note that the equation (7) has a unique solution, because  $\alpha_{\mu}$  and  $\beta_{\mu}$ , considered as functions of  $\mu > 0$ , are non-negative, continuous, non-increasing and tend to zero as  $\mu \to +\infty$ .

**3.3. Proof of Theorem 1.3.2.** Consider some potential  $q^* \in -\Gamma_1$ , and some positive eigenfunction  $y \in \ker T_{q^*}(\lambda_1(q^*))$ . Suppose that  $\operatorname{supp} q^* \subseteq \Omega^-(y)$ . Then for any  $q \in -\Gamma_1$  we, using 2.3.3, have

$$\begin{aligned} 0 &= \left\langle T_{q^*} \left( \lambda_1(q^*) \right) y, y \right\rangle \\ &= \int_0^1 \left[ (y')^2 - \lambda_1(q^*) y^2 \right] dx - \inf_{x \in [0,1]} y^2(x) + k_0^2 y^2(0) + k_1^2 y^2(1) \\ &\geq \int_0^1 \left[ (y')^2 + \left( q - \lambda_1(q^*) \right) y^2 \right] dx + k_0^2 y^2(0) + k_1^2 y^2(1), \end{aligned}$$

hence

$$\left\langle T_q(\lambda_1(q^*))y, y \right\rangle \le 0.$$

It follows that  $\lambda_1(q) \leq \lambda_1(q^*)$ , therefore  $\lambda_1(q^*) = M_1^-$ . Thus we have proved that  $M_1^-$  is attained at any potential  $q^*$  such that  $\sup q^* \subseteq \Omega^-(y)$ .

Suppose  $k_0^2 + k_1^2 \leq 1$ . Consider the generalized function

$$q^* \rightleftharpoons -k_0^2 \boldsymbol{\delta}_0 - k_1^2 \boldsymbol{\delta}_1 - \left(1 - k_0^2 - k_1^2\right),$$

which in this case belongs to  $-\Gamma_1$ . Using (14), we get that the first eigenfunction of the pencil  $T_{q^*}$  is  $y \equiv \text{const}$ , so  $\text{supp } q^* \subseteq \Omega^-(y)$ . It follows that  $M_1^-$  is attained at the potential  $q^*$  and is equal to the corresponding first eigenvalue

$$\lambda_1(q^*) = k_0^2 + k_1^2 - 1.$$

Suppose

$$k_0^2 + k_1^2 \ge 1,\tag{35}$$

$$k_1^2 - k_0^2 \le 1. \tag{36}$$

Consider the generalized function  $q^* \rightleftharpoons -(1+k_0^2-k_1^2)\boldsymbol{\delta}_0/2 - (1-k_0^2+k_1^2)\boldsymbol{\delta}_1/2$ , which, due to (36), belongs to  $-\Gamma_1$ . For such  $q^*$  the equation  $T_{q^*}(\lambda)y = 0$  is equivalent to the problem (8), (9). The first eigenvalue  $\lambda_1(q^*)$ , due to (35) and (9), is non-negative and the corresponding eigenfunction is

$$y(x) \equiv \cos\left[\sqrt{\lambda_1(q^*)} \cdot (x-\zeta)\right],\tag{37}$$

where  $\zeta = 1/2$ . Hence supp  $q^* \subseteq \Omega^-(y)$ . It follows that  $M_1^-$  is attained at the potential  $q^*$  and is equal to the corresponding first eigenvalue  $\lambda_1(q^*)$ .

Suppose  $k_1^2 - k_0^2 \ge 1$ . Consider the generalized function  $q^* \rightleftharpoons -\delta_1 \in -\Gamma_1$ . For such  $q^*$  the equation  $T_{q^*}(\lambda)y = 0$  is equivalent to the problem (8), (10).

The corresponding first eigenfunction is defined by (37), where  $\zeta \in [0, 1/2]$ , since  $k_1^2 - 1 \ge k_0^2$ . Hence supp  $q^* \subseteq \Omega^-(y)$ . It follows that  $M_1^-$  is attained at the potential  $q^*$  and is equal to the corresponding first eigenvalue  $\lambda_1(q^*)$ .

**3.4. Proof of Theorem 1.3.3.** Consider some potential  $q \in \Gamma_1$ , and some positive eigenfunction  $y \in \ker T_q(\lambda_1(q))$ . Then for any  $\lambda > \lambda_1(q)$ , according to 2.3.3, we have

$$0 > \int_{0}^{1} \left[ (y')^{2} + (q - \lambda) y^{2} \right] dx + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1)$$
  
$$\geq \int_{0}^{1} \left[ (y')^{2} - \lambda y^{2} \right] dx + \inf_{x \in [0,1]} y^{2}(x) + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1).$$

It follows that there exists  $\zeta \in [0, 1]$  such that

$$\int_{0}^{1} \left[ (y')^{2} + (\boldsymbol{\delta}_{\zeta} - \lambda) y^{2} \right] dx + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1) < 0.$$

So for any  $\lambda > m_1^+$  there exists  $\zeta \in [0,1]$  such that  $\lambda_1(\boldsymbol{\delta}_{\zeta}) < \lambda$ . Hence, using 2.3.3, we get  $m_1^+ = \inf_{x \in [0,1]} \lambda_1(\boldsymbol{\delta}_x)$ . This equality is equivalent, according to 2.4.1, to the following fact:  $F(m_1^+, x)$  is defined for all  $x \in [0,1]$  and satisfies  $\sup_{x \in [0,1]} F(m_1^+, x) = 1$ .

Since  $m_1^+ > 0$ , from 2.4.3 it follows that if  $\mu = m_1^+$ , then for any  $\zeta \in [0, 1]$  the conditions (25) hold. According to (26), (25) and

$$\frac{\partial F(\mu,\zeta)}{\partial \zeta} \equiv \mu \cdot \frac{\cos^2[\sqrt{\mu} \cdot (1-\beta_{\mu}-\zeta)] - \cos^2[\sqrt{\mu} \cdot (\zeta-\alpha_{\mu})]}{\cos^2[\sqrt{\mu} \cdot (\zeta-\alpha_{\mu})] \cdot \cos^2[\sqrt{\mu} \cdot (1-\beta_{\mu}-\zeta)]},\tag{38}$$

it follows that the function  $F(\mu, \cdot)$  can have at some point  $\zeta \in (0, 1)$  a local extremum satisfying  $F(\mu, \zeta) > 0$  only if  $\zeta = (1 - \beta_{\mu} + \alpha_{\mu})/2$ ,  $\zeta > \alpha_{\mu}$  and  $\zeta < 1 - \beta_{\mu}$ . But this conditions imply, according to (38), that such  $\zeta$  must be a point of strict local minimum of the function  $F(\mu, \cdot)$ . Therefore,  $F(\mu, \cdot)$  cannot have a supremum in (0, 1), so we get  $m_1^+ = \inf\{\lambda_1(\delta_0), \lambda_1(\delta_1)\}$ . Note that for the potential  $q^* \rightleftharpoons \delta_i$ , where  $i \in \{0, 1\}$ , the equation  $T_{q^*}(\lambda)y = 0$  is equivalent to the problem

$$-y'' = \lambda y,$$
  
$$y'(0) - \left[k_0^2 + (1-i)\right] y(0) = y'(1) + \left[k_1^2 + i\right] y(1) = 0.$$

Therefore, we have

$$\frac{\lambda_1(\boldsymbol{\delta}_i) - k_0^2 k_1^2 - k_{1-i}^2}{k_0^2 + k_1^2 + 1} = \sqrt{\lambda_1(\boldsymbol{\delta}_i)} \cot \sqrt{\lambda_1(\boldsymbol{\delta}_i)},$$

so  $m_1^+ = \lambda_1(\boldsymbol{\delta}_1).$ 

**3.5. Proof of Theorem 1.3.4.** Consider some potential  $q \in -\Gamma_1$ , and some positive eigenfunction  $y \in \ker T_q(\lambda_1(q))$ . Then for any  $\lambda > \lambda_1(q)$ , according to 2.3.3, we have

$$\begin{split} 0 > & \int_{0}^{1} \left[ (y')^{2} + (q - \lambda) y^{2} \right] dx + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1) \\ \ge & \int_{0}^{1} \left[ (y')^{2} - \lambda y^{2} \right] dx - \sup_{x \in [0,1]} y^{2}(x) + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1). \end{split}$$

It follows that there exists  $\zeta \in [0, 1]$  such that

$$\int_{0}^{1} \left[ (y')^{2} + (-\delta_{\zeta} - \lambda) y^{2} \right] dx + k_{0}^{2} y^{2}(0) + k_{1}^{2} y^{2}(1) < 0.$$

So for any  $\lambda > m_1^-$  there exists  $\zeta \in [0, 1]$  such that  $\lambda_1(-\delta_{\zeta}) < \lambda$ . Hence, using 2.3.3, we get  $m_1^- = \inf_{x \in [0,1]} \lambda_1(-\delta_x)$ . This equality is equivalent, according to 2.4.1, to the following fact:  $F(m_1^-, x)$  is defined for all  $x \in [0, 1]$  and satisfies  $\sup_{x \in [0,1]} F(m_1^-, x) = -1$ .

For any fixed value  $\mu \in \mathbb{R}$  we consider the conditions

$$F(\mu,\zeta) < 0, \tag{39}$$

$$\partial F(\mu,\zeta)/\partial\zeta = 0.$$
 (40)

It is clear that some point  $\zeta \in (0,1)$  can satisfy the equalities  $F(\mu,\zeta) = \sup_{x \in [0,1]} F(\mu,x) = -1$  only if (39) and (40) hold.

Suppose  $\mu > 0$ . Then, according to (38), (26), (32) and (30), for any point  $\zeta \in (0, 1)$  satisfying (39) the condition (40) holds if and only if the problem

$$-y'' = \mu y$$
 at  $(0, \zeta) \cup (\zeta, 1),$  (41)

$$y'(0) - k_0^2 y(0) = 2y'(\zeta - 0) + F(\mu, \zeta)y(\zeta)$$
  
= 2y'(\zeta + 0) - F(\mu, \zeta)y(\zeta) = y'(1) + k\_1^2 y(1) = 0 (42)

has a continuous positive solution. Besides, for any point  $\zeta \in (0, 1)$  satisfying (39) and (40) we have

$$\alpha_{\mu} > \zeta > 1 - \beta_{\mu}. \tag{43}$$

Therefore, according to (38) and (25), this stationary point  $\zeta$  is a strict maximum of  $F(\mu, \cdot)$ . Since for any  $x \in [0, 1]$ , using (43), we get

$$-\pi/2 < -\sqrt{\mu}\alpha_{\mu} \le \sqrt{\mu} \cdot (x - \alpha_{\mu}) < \sqrt{\mu} \cdot (x - 1 + \beta_{\mu}) \le \sqrt{\mu}\beta_{\mu} < \pi/2,$$

it follows from the proposition 2.4.3 that the function  $F(\mu, \cdot)$  is defined everywhere on [0, 1].

Suppose  $\mu = 0$ . Let us use the same method as in the previous case, changing (26) to (27), and (32) to (33). Then we get that for any point  $\zeta \in (0, 1)$  satisfying (39) the condition (40) holds if and only if the problem (41), (42) has a continuous positive solution. Using (27) we also get that the second derivative of  $F(0, \cdot)$  is negative. Hence any stationary point  $\zeta \in (0, 1)$  is a strict maximum of  $F(0, \cdot)$ .

Suppose  $\mu \in (-k_0^4, 0)$ . Then, using (28), we get

$$\frac{\partial F(\mu,\zeta)}{\partial \zeta} \equiv -\mu \bigg\{ \sinh^{-2} \left( \sqrt{|\mu|} \zeta + \alpha_{\mu} \right) - \sinh^{-2} \left( \sqrt{|\mu|} (1-\zeta) + \beta_{\mu} \right) \bigg\},\,$$

where

$$\alpha_{\mu} \rightleftharpoons \frac{1}{2} \ln \frac{k_0^2 + \sqrt{|\mu|}}{k_0^2 - \sqrt{|\mu|}}, \quad \beta_{\mu} \rightleftharpoons \frac{1}{2} \ln \frac{k_1^2 + \sqrt{|\mu|}}{k_1^2 - \sqrt{|\mu|}}$$

Therefore, according to (34), for any point  $\zeta \in (0, 1)$  satisfying (39) the condition (40) holds if and only if the problem (41), (42) has a continuous positive solution. Since  $\partial^2 F(\mu, \zeta)/\partial \zeta^2 < 0$ , it follows that any stationary point  $\zeta \in (0, 1)$ is a strict maximum of  $F(\mu, \cdot)$ .

Suppose  $0 > \mu = -k_0^4 = -k_1^4$ . Then the function  $F(\mu, \cdot)$  is a negative constant, and for any point  $\zeta \in (0, 1)$  problem (41), (42) has a continuous positive solution.

Suppose  $\mu \in [-k_1^4, -k_0^4]$ , also  $\mu < 0$  and  $k_1 > k_0$ . Then from (28) and (34) it follows that  $\partial F(\mu, \zeta) / \partial \zeta < 0$ , and the problem (41), (42) has no positive solutions for any  $\zeta \in (0, 1)$ .

Suppose  $\mu < -k_1^4$ . Then, using (28), we get

$$\frac{\partial F(\mu,\zeta)}{\partial \zeta} \equiv \mu \bigg\{ \cosh^{-2} \left( \sqrt{|\mu|} \zeta + \alpha_{\mu} \right) - \cosh^{-2} \left( \sqrt{|\mu|} (1-\zeta) + \beta_{\mu} \right) \bigg\},\$$

where

$$\alpha_{\mu} \rightleftharpoons \frac{1}{2} \ln \frac{\sqrt{|\mu|} + k_0^2}{\sqrt{|\mu|} - k_0^2}, \quad \beta_{\mu} \rightleftharpoons \frac{1}{2} \ln \frac{\sqrt{|\mu|} + k_1^2}{\sqrt{|\mu|} - k_1^2}$$

Therefore, according to (34), for any point  $\zeta \in (0, 1)$  satisfying (39) the condition (40) holds if and only if problem (41), (42) has a continuous positive solution. Since  $\partial^2 F(\mu, \zeta)/\partial \zeta^2 > 0$ , it follows that any stationary point  $\zeta \in (0, 1)$  is a strict minimum of  $F(\mu, \cdot)$ .

From the proposition 2.4.1 we also get that for any  $\mu \leq 0$  the function  $F(\mu, \cdot)$  is defined everywhere on [0, 1].

Combining all this, we obtain the following: the existence of a continuous positive solution to the problem (12), (13) for some  $\mu \geq -k_0^4$  and  $\zeta \in (0,1)$  implies that  $F(\mu, \zeta) = -1$ , the function  $F(\mu, \cdot)$  is defined everywhere on [0, 1], and  $\sup_{x \in [0,1]} F(\mu, x) \leq -1$ . Therefore,  $m_1^- = \lambda_1(-\boldsymbol{\delta}_{\zeta})$ . In converse, if for any  $\mu \geq -k_0^4$  and  $\zeta \in (0,1)$  the positive solution of (12), (13) does not exist, we get  $m_1^- = \inf\{\lambda_1(-\boldsymbol{\delta}_0), \lambda_1(-\boldsymbol{\delta}_1)\}$ . From the equation

$$\lambda_1(-\boldsymbol{\delta}_i) - k_0^2 k_1^2 + k_{1-i}^2 = (k_0^2 + k_1^2 - 1) \cdot \psi(\lambda_1(-\boldsymbol{\delta}_i)),$$

where  $i \in \{0, 1\}$  and

$$\psi(x) \rightleftharpoons \begin{cases} \sqrt{x} \cot \sqrt{x} & \text{for } x > 0, \\ 1 & \text{for } x = 0, \\ \sqrt{|x|} \coth \sqrt{|x|} & \text{for } x < 0, \end{cases}$$

we obtain that  $\inf \{\lambda_1(-\boldsymbol{\delta}_0), \lambda_1(-\boldsymbol{\delta}_1)\} = \lambda_1(-\boldsymbol{\delta}_0).$ 

**3.6.** Now we get some conditions for the existence of a continuous positive solution to the problem (12), (13) considered in Theorem 1.3.4.

Suppose  $\mu_0(\zeta)$ , where  $\zeta \in (0, 1]$ , is the minimal eigenvalue of the problem

$$-y'' = \lambda y, \tag{44}$$

$$y'(0) - k_0^2 y(0) = 2y'(\zeta) - y(\zeta) = 0,$$
(45)

and suppose  $\mu_1(\zeta)$ , where  $\zeta \in [0, 1)$ , is the minimal eigenvalue of the problem (44) and

$$2y'(\zeta) + y(\zeta) = y'(1) + k_1^2 y(1) = 0.$$

It is clear that for some  $\mu \in \mathbb{R}$  and  $\zeta \in (0, 1)$  a continuous positive solution to (12), (13) exists if and only if the equalities  $\mu_0(\zeta) = \mu_1(\zeta) = \mu$  hold.

**3.6.1.** If  $k_0^2 = 1/2$ , then  $\mu_0(\zeta) \equiv -1/4$ . If  $k_0^2 > 1/2$ , then the function  $\mu_0$  strictly decreases and satisfies

$$\lim_{\zeta \to 0} \mu_0(\zeta) = +\infty \quad and \quad \mu_0(1) > -1/4.$$

If  $k_0^2 < 1/2$ , then for any  $\zeta \in (0,1]$  the inequality  $\mu_0(\zeta) < -1/4$  holds.

Proof. Suppose  $k_0^2 = 1/2$ . Then for any  $\zeta \in (0, 1]$  the problem (44), (45) has the positive eigenfunction  $y(x) \equiv e^{x/2}$  corresponding to the eigenvalue -1/4.

Suppose  $k_0^2 > 1/2$ . Since the eigenvalues of the problem (44), (45) increase by  $k_0^2$ , it follows that  $\mu_0(\zeta) > -1/4$ . Then let  $y_0 \in W_2^1[0, \zeta]$  be an eigenfunction of the problem (44), (45) corresponding to the eigenvalue  $\mu_0(\zeta)$ . Continuing the function  $y_0$  for any  $\theta \in (\zeta, 1]$  to the interval  $(\zeta, \theta]$  in the form  $y(x) \rightleftharpoons y_0(\zeta)e^{(x-\zeta)/2}$ , for the obtained function  $y \in W_2^1[0, \theta]$  we get

$$\int_{0}^{\theta} [(y')^{2} - \mu_{0}(\zeta) y^{2}] dx + k_{0}^{2} y^{2}(0) - \frac{y^{2}(\theta)}{2} = [-1/4 - \mu_{0}(\zeta)] \cdot [e^{\theta - \zeta} - 1] \cdot y^{2}(\zeta) < 0,$$

hence  $\mu_0(\theta) < \mu_0(\zeta)$ .

Finally, for  $\zeta \to 0$  we have uniform by  $y \in W_2^1[0, \zeta]$  asymptotic estimate

$$\begin{split} &\int_{0}^{\zeta} (y')^2 \, dx + k_0^2 \, y^2(0) - \frac{y^2(\zeta)}{2} \\ &= \left[ \int_{0}^{\zeta} \frac{(y')^2}{2} \, dx + (k_0^2 - 1/2) \, y^2(0) \right] + \left[ \int_{0}^{\zeta} \frac{(y')^2}{2} \, dx + \frac{y^2(0) - y^2(\zeta)}{2} \right] \\ &\geq \frac{k_0^2 - 1/2 + o(1)}{\zeta} \int_{0}^{\zeta} y^2 \, dx - \frac{1}{2} \int_{0}^{\zeta} y^2 \, dx, \end{split}$$

therefore,  $\mu_0(\zeta) \ge [k_0^2 - 1/2 + o(1)] \cdot \zeta^{-1}$ .

The inequality  $\mu_0(\zeta) < -1/4$  for the case  $k_0^2 < 1/2$  is proved likewise the inequality  $\mu_0(\zeta) > -1/4$  for the case  $k_0^2 > 1/2$ .

**3.6.2.** If  $k_1^2 = 1/2$ , then  $\mu_1(\zeta) \equiv -1/4$ . If  $k_1^2 > 1/2$ , then the function  $\mu_1$  strictly increases and satisfies

$$\lim_{\zeta \to 1} \mu_1(\zeta) = +\infty \quad and \quad \mu_1(0) > -1/4.$$

If  $k_1^2 < 1/2$ , then for any  $\zeta \in [0,1)$  the inequality  $\mu_1(\zeta) < -1/4$  holds.

The proposition 3.6.2 is proved likewise 3.6.1.

Combining 3.6.1 and 3.6.2, we get the last proposition:

**3.6.3.** The problem (12), (13) has a continuous positive solution for some  $\mu \geq -k_0^4$  and  $\zeta \in (0,1)$  if and only if this condition holds:

$$k_0^2 > 1/2$$
 or  $k_0^2 = k_1^2 = 1/2$ .

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